Medical Imaging

Prof. Dr. Tobias Knopp October 25, 2022

Institute für Biomedizinische Bildgebung

Signal Processing

- Understanding tomographic imaging techniques requires profound knowledge of signal processing
- In particular Fourier analysis plays an important role
- Basics of signal processing will be recapitulated

- A signal can be seen as a message that is send from a sender to a recipient.
- Usually they transport a physical quantity
- Signals can either be discrete or continuous (i.e. a function of ${\mathbb N}$ or ${\mathbb R})$
- Typically a signal depends on *time* and/or *space*

Signal Examples

- s(x): spatial 1D signal, e.g. detector array in CT
- f(x, y): spatial 2D signal, e.g. slice through an object in CT
- c(t): temporal 1D signal, e.g. ECG



Fundamental Signals

• Heaviside step function

$$\operatorname{step}(x) = \begin{cases} 1, & ext{for } x \geq 0 \\ 0, & ext{for } x < 0 \end{cases}$$

• Rectangular function

$$\operatorname{rect}(x) = \begin{cases} 1, & \text{ for } |x| < \frac{1}{2} \\ 0.5, & \text{ for } |x| = \frac{1}{2} \\ 0, & \text{ for } |x| > \frac{1}{2} \end{cases}$$

• Sinc function

$$\operatorname{si}(x) = \operatorname{sinc}(x) = \frac{\operatorname{sin}(\pi x)}{\pi x}$$

(3)

• Dirac delta distribution

$$\delta(x - x_0) = \begin{cases} 0, & \text{for } x \neq x_0 \\ \infty, & \text{for } x = x_0 \end{cases}$$
$$= \lim_{\tilde{x} \to 0} \frac{1}{\tilde{x}} \operatorname{rect}(\frac{x - x_0}{\tilde{x}})$$

- The Dirac delta distribution is no function in a classical sense as ∞ can only be considered in the limes and is not a valid function value.
- The Dirac delta distribution can be defined over the integral

$$\int_{-\infty}^{\infty} \delta(x - x_0) f(x) \, \mathrm{d}x = f(x_0) \tag{4}$$

In particular it holds that

$$\int_{-\infty}^{\infty} \delta(x - x_0) \, \mathrm{d}x = 1 \tag{5}$$

I.e. Dirac delta distribution has unit area.

Systems

• A system \mathcal{L} takes as input a function f(x) and outputs a function g(x)

 $g(x) = \mathcal{L}\{f(x)\}$

• A system is linear if

$$\mathcal{L}\{\sum_{i}a_{i}f_{i}(x)\}=\sum_{i}a_{i}\mathcal{L}\{f_{i}(x)\}$$

• A system is said to be time invariant / shift invariant if

$$\mathcal{L}{f(x-x_0)} = \mathcal{L}{f(x-x_0)}.$$

This means a shift of the input signal by x_0 leads to a shift by x_0 of the output signal.

• A system is said to be LTI or LSI if it is linear and time / shift invariant.

Impulse Response and Convolution

Any LSI system can be described by a convolution integral

$$g(x) = \mathcal{L}{f(x)}$$
$$= \int_{-\infty}^{\infty} f(\tilde{x})h(x - \tilde{x}) d\tilde{x}$$
$$= (f * h)(x)$$

h is the so-called *impulse response* or *point-spread function (PSF)* that is obtained by applying a Dirac delta to the LSI system

$$\mathcal{L}\{\delta(x)\} = \int_{-\infty}^{\infty} \delta(\tilde{x})h(x - \tilde{x}) d\tilde{x}$$

=
$$\int_{-\infty}^{\infty} \delta(x - y)h(y) dy$$
$$= \int_{-\infty}^{\infty} \delta(y - x)h(y) dy = h(x)$$

Derivation of the Convolution Integral

Idea: express f(x) as a sum of shifted rectangular functions

$$f(x) = \lim_{x_0 \to 0} \sum_{n = -\infty}^{\infty} f(nX_0) \operatorname{rect}\left(\frac{x - nX_0}{X_0}\right)$$

Applying LSI properties yields

$$g(x) = \mathcal{L}{f(x)} = \lim_{X_0 \to 0} \sum_{n = -\infty}^{\infty} f(nX_0)\mathcal{L}\left\{\operatorname{rect}\left(\frac{x - nX_0}{X_0}\right)\right\}$$
$$= \int_{-\infty}^{\infty} f(\tilde{x})\mathcal{L}{\delta(x - \tilde{x})} d\tilde{x}$$
$$= \int_{-\infty}^{\infty} f(\tilde{x})h(x - \tilde{x}) d\tilde{x}$$

It is known from analysis lectures that (almost) any *p*-periodic function s(x) can be expanded into a *Fourier series*

$$s(x) = \sum_{n=-\infty}^{\infty} c_n \mathrm{e}^{2\pi \mathrm{i} \frac{nx}{p}}$$

consisting of complex sinus functions $e^{ix} = \cos x + i \sin x$. The Fourier coefficients c_n can be calculated by

$$c_n = \frac{1}{p} \int_{-p/2}^{p/2} s(x) \mathrm{e}^{-2\pi \mathrm{i} \frac{nx}{p}} \,\mathrm{d}x$$

In order to also express non-periodic functions by Fourier series (i.e. any general signal/function) one can consider the limes $p \to \infty$

$$s(x) = \lim_{p \to \infty} \sum_{n = -\infty}^{\infty} c_n e^{2\pi i \frac{nx}{p}}$$
$$= \int_{f:=\frac{n}{p}}^{\infty} \int_{-\infty}^{\infty} S(f) e^{2\pi i fx} df =: \mathcal{F}^{-1}\{S(f)\}$$

Here, $f := \frac{n}{p}$ is the frequency and $S(f) = S(\frac{n}{p}) = c_n$ is a continuous function of frequency. This is in contrast to the discrete spectrum of the periodic Fourier series.

The Fourier transform S(f) of s(x) can be calculated by

$$S(f) = \int_{-\infty}^{\infty} s(x) \mathrm{e}^{-2\pi \mathrm{i} f x} \, \mathrm{d} x =: \mathcal{F}\{s(x)\}$$

S(f) is often called the *spectrum* of s(x). The Fourier relation is often indicated by

$$s(x) \longrightarrow S(f)$$

Example

The Fourier transformation of the rectangular function can be calculated to be

$$\mathcal{F}\{\operatorname{rect}(x)\} = \int_{-\infty}^{\infty} \operatorname{rect}(x) e^{-2\pi i f x} dx$$
$$= \int_{-\frac{1}{2}}^{\frac{1}{2}} e^{-2\pi i f x} dx$$
$$= \frac{i}{2\pi f} \left(e^{-i\pi f} - e^{i\pi f} \right)$$
$$= \frac{\sin(\pi f)}{\pi f} = \operatorname{sinc}(f)$$

In the last step the euler formula $sin(f) = \frac{e^{if} - e^{-if}}{2i}$ has been used.

Instead of describing an LSI system as a convolution with the PSF h(x) one can equivalently describe it in Fourier space by considering the *transfer function* $H(f) = \mathcal{F}{h(x)}$.

Theorem

A convolution g(x) = (s * h)(x) in spatial space corresponds to a multiplication in Fourier space:

G(f)=S(f)H(f)

where $G(f) = \mathcal{F}{g(x)}$, $S(f) = \mathcal{F}{s(x)}$, and $H(f) = \mathcal{F}{h(x)}$.

Proof

$$G(f) = \mathcal{F}\{g(x)\} = \mathcal{F}\{(s * h)(x)\}$$

= $\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} s(\tilde{x})h(x - \tilde{x}) d\tilde{x}e^{-2\pi i f x} dx$
= $\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} s(\tilde{x})h(x - \tilde{x})e^{-2\pi i f(x - \tilde{x})}e^{-2\pi i f \tilde{x}} d\tilde{x} dx$
= $\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} s(\tilde{x})h(z)e^{-2\pi i f \tilde{x}}e^{-2\pi i f z} d\tilde{x} dz$
= $\int_{-\infty}^{\infty} s(\tilde{x})e^{-2\pi i f \tilde{x}} d\tilde{x} \int_{-\infty}^{\infty} h(z)e^{-2\pi i f z} dz$
= $S(f)H(f)$

- Efficient application of filter (low pass, high pass)
- Deconvolution / image sharpening

In order to numerically calculate Fourier coefficients one has to sample the Fourier integral a discrete and equidistant sampling points.

This makes the Fourier coefficients periodic so that the data in both domains (spatial and frequency) is discrete (line spectrum) and periodic.

Given a sequence s_0, \ldots, s_{N-1} the discrete Fourier transformation (DFT) is defines as

$$S_m = \sum_{n=0}^{N-1} s_n \mathrm{e}^{-2\pi \mathrm{i} \frac{nm}{N}}, \quad m = 0, \dots, N-1$$

The DFT is a unitary transformation and its inverse can be calculated by

$$s_n = \frac{1}{N} \sum_{m=0}^{N-1} S_m \mathrm{e}^{2\pi \mathrm{i} \frac{nm}{N}}, \quad n = 0, \dots, N-1$$

When defining the vectors $\boldsymbol{S} := (S_m)_{m=0}^{N-1}$ and $\boldsymbol{s} := (s_n)_{n=0}^{N-1}$, and the discrete Fourier matrix $\boldsymbol{F} := \left(e^{-2\pi i \frac{nm}{N}}\right)_{m=0,\ldots,N-1;n=0,\ldots,N-1}$ the DFT can be written in matrix-vector form

S = Fs

A naive implementation of the DFT would require $\mathcal{O}(N^2)$ arithmetic operations.

The fast Fourier transformation is a fast algorithm capable of carrying out FFT in only $\mathcal{O}(N \log N)$. It uses a recursive divide and conquer principle.

Important is that the FFT can only be applied to equidistant sampling positions.

Assumption: We will derive the DFT for $N = 2^r$.

Basic idea: split the sum into two sums, which are itself regular DFTs.

$$S_{m} = \sum_{n=0}^{N-1} s_{n} e^{-2\pi i \frac{nm}{N}}$$
$$= \sum_{n=0}^{\frac{N}{2}-1} s_{n} e^{-2\pi i \frac{nm}{N}} + \sum_{n=0}^{\frac{N}{2}-1} s_{n+\frac{N}{2}} e^{-2\pi i \frac{(n+\frac{N}{2})m}{N}}$$

Now we will discuss two cases. m = 2l (even) and m = 2l + 1 (odd) for $l = 0, ..., \frac{N}{2} - 1$.

Fast Fourier Transformation

Case m = 2/:

$$S_{2l} = \sum_{n=0}^{\frac{N}{2}-1} s_n e^{-2\pi i \frac{n2l}{N}} + \sum_{n=0}^{\frac{N}{2}-1} s_{n+\frac{N}{2}} e^{-2\pi i \frac{(n+\frac{N}{2})2l}{N}}$$
$$= \sum_{n=0}^{\frac{N}{2}-1} s_n e^{-2\pi i \frac{nl}{N}} + \sum_{n=0}^{\frac{N}{2}-1} s_{n+\frac{N}{2}} e^{-2\pi i \frac{nl}{N}} \underbrace{e^{-2\pi i l}}_{1}$$
$$= \underbrace{\sum_{n=0}^{\frac{N}{2}-1} (s_n + s_{n+\frac{N}{2}}) e^{-2\pi i \frac{nl}{N}}}_{\text{DFT length } \frac{N}{2}}$$

Fast Fourier Transformation

Case m = 2l + 1: $S_{2l+1} = \sum_{n=0}^{\frac{N}{2}-1} s_n e^{-2\pi i \frac{n(2l+1)}{N}} + \sum_{n=0}^{\frac{N}{2}-1} s_{n+\frac{N}{2}} e^{-2\pi i \frac{(n+\frac{N}{2})(2l+1)}{N}}$ $=\sum_{n=0}^{\frac{N}{2}-1} s_{n} e^{-2\pi i \frac{n!}{2}} e^{-2\pi i \frac{n}{N}} + \sum_{n=0}^{\frac{N}{2}-1} s_{n+\frac{N}{2}} e^{-2\pi i \frac{n!}{2}} \underbrace{e^{-2\pi i \frac{n}{N}}}_{1} e^{-2\pi i \frac{n}{N}} \underbrace{e^{-2\pi i \frac{n}{2}}}_{-1}$ $=\sum_{j=1}^{\frac{N}{2}-1} s_{n} \mathrm{e}^{-2\pi \mathrm{i}\frac{nj}{N}} \mathrm{e}^{-2\pi \mathrm{i}\frac{n}{N}} - \sum_{j=1}^{\frac{N}{2}-1} s_{n+\frac{N}{2}} \mathrm{e}^{-2\pi \mathrm{i}\frac{nj}{N}} \mathrm{e}^{-2\pi \mathrm{i}\frac{n}{N}}$ $\frac{N}{2} - 1$ $=\sum_{n=0}^{\infty} e^{-2\pi i \frac{n}{N}} (s_n - s_{n+\frac{N}{2}}) e^{-2\pi i \frac{n!}{N}}$ DFT length $\frac{N}{2}$

22

Instead of one length N DFT we can apply two length $\frac{N}{2}$ DFTs.

The transformation can be applied again and in the second step we need to apply four length $\frac{N}{4}$ DFTs.

After $r = \log(N)$ steps we need to apply N DFTs of length 1.

In each step the algorithm needs $\frac{N}{2}$ additions, $\frac{N}{2}$ subtractions, and $\frac{N}{2}$ multiplications. The algorithmic complexity it thus $\mathcal{O}(N)$ arithmetic operations.

Since the DFT requires $\log(N)$ steps, the overal time complexity is $\mathcal{O}(N \log(N))$ compared to $\mathcal{O}(N^2)$ of an ordinary DFT.

Fast Fourier Transformation

The FFT is usually implemented inplace and can be visualized as follows



- We have derived the FFT for $N = 2^r$
- The same can be done for other basis 3,5,7,11,13, ...
- By prime factorization one can apply the FFT to general N
- Alternatively, one can pad the vector with zeros to the next power of 2
- The most popular FFT library is the FFTW (Fastest Fourier Transform in the West). It is used in the Julia package FFTW.jl